

# QCD on the lattice - an introduction

## Lecture 5

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# Baryon correlation functions

- Baryons have two-point function representations too. The simplest creation operator would be







$$\Phi(\underline{x}, t) = \Gamma_{\alpha\beta\gamma} \epsilon_{ijk} u_i^\alpha(\underline{x}, t) u_j^\beta(\underline{x}, t) d_k^\gamma(\underline{x}, t)$$

$\Gamma$  is the spin contraction to make irreducible representations of the double cover of  $O_h$ .

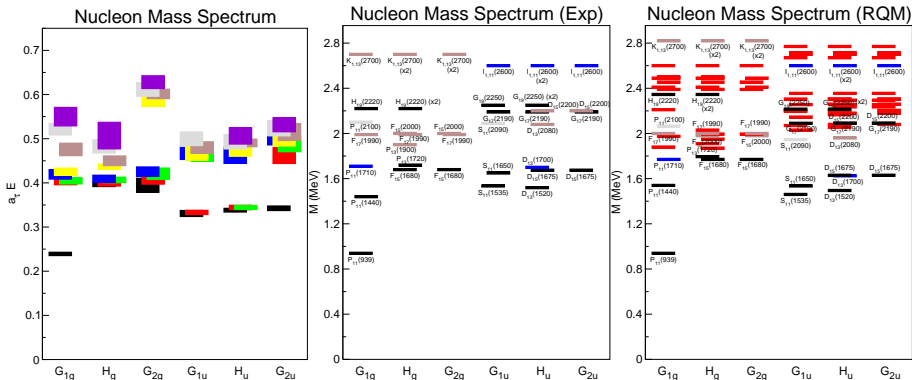
- There are six double cover representations:  $G_1^{u,g}$   $G_2^{u,g}$   $H^{u,g}$
- More complicated operators can be constructed again using gluon links to make gauge invariants.

# Baryon correlation functions

- Example from [Phys.Rev.D72:094506,2005](#).

Operator type	Displacement indices
 single-site	$i = j = k = 0$
 singly-displaced	$i = j = 0, k \neq 0$
 doubly-displaced-I	$i = 0, j = -k, k \neq 0$
 doubly-displaced-L	$i = 0,  j  \neq  k , jk \neq 0$
 triply-displaced-T	$i = -j,  j  \neq  k , jk \neq 0$
 triply-displaced-O	$ i  \neq  j  \neq  k , ijk \neq 0$

# Baryon spectroscopy



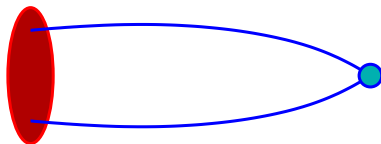
From thesis of Adam Lichtl (CMU)

# Matrix elements

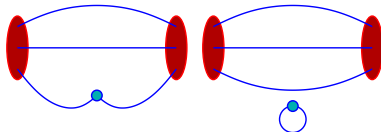
- Physics including the interface of QCD with the other fields in the standard model can be investigated on the lattice.
- Examples:
  - 1 The pion decay constant,  $f_\pi$
  - 2 The B-meson decay constant,  $f_B$
  - 3 Bag parameters  $B_K$
  - 4 Form factors for electromagnetic probes interacting with mesons and baryons
- These require the evaluation of QCD matrix elements, and usually involve renormalisation to link to the other SM fields.
- Perturbation theory converges if the matching scale is high and there is a lot of work on non-perturbative matching methods.

# Matrix elements

- $f_\pi$  requires the evaluation of a two-point function:



- Here, the blue dot represents the insertion of the axial vector current operator. Linking with the rest of the SM requires renormalisation/matching.
- Form factors require the evaluation of diagrams like:



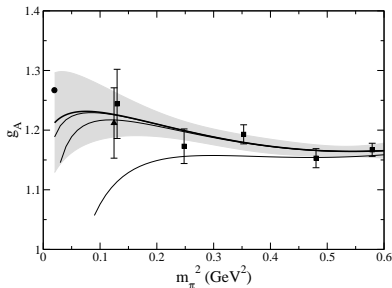
# Example of current matrix-element calculation

- The axial vector  $A_\mu = \bar{q}\gamma_\mu\gamma_5\tau q$  form factor of the nucleon (whose size governs  $\beta$  decay) is

$$\langle N(p+q, t') | A_\mu(q) | N(p, t) \rangle = \bar{u}(p+q, t') \frac{\tau}{2} [g_A(q^2)\gamma_\mu\gamma_5 + g_P(q^2)q_\mu\gamma_5] u(p, t)$$

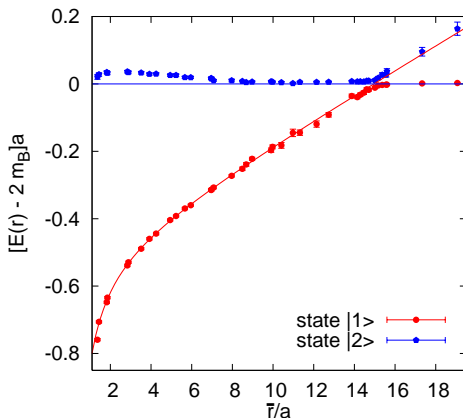
LHPC collaboration - Phys.Rev.Lett. 96 (2006) 052001

- Nucleon axial charge  $g_A(0)$  vs pion mass. Experimental result is the black circle (left).



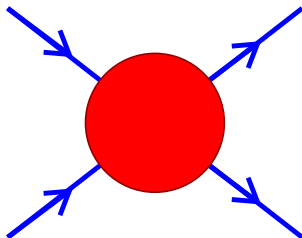
# Decay physics: breaking the QCD string

- With fundamental representation quarks, the gluon string that generates the confining potential can be broken by vacuum creation of a quark anti-quark pair
- This physics is missing in the Yang-Mills theory, but present in QCD.
- This has been observed on the lattice. [Phys.Rev.D71:114513,2005](#).





# The bad news - Maiani-Testa No-go theorem



- Importance sampling Monte Carlo simulation only works efficiently for a path integral with a positive definite probability measure: Euclidean space.
- Maiani-Testa: Scattering matrix elements cannot be extracted from infinite-volume Euclidean-space correlation functions (except at threshold).
- Can the lattice tell us anything about low-energy scattering or states above thresholds?

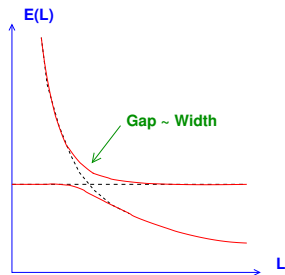
# Scattering lengths indirectly: Lüscher's method

- Scattering lengths can be inferred indirectly given the right measurements in Euclidean field theory.
- In a three-dimensional box with finite size  $L$ , the spectrum of low-lying states is discrete, even above thresholds (since the momenta of daughter mesons are quantised).
- Precise data on the dependence of the energy spectrum on  $L$  can be used to compute low-energy scattering (below inelastic threshold).
- This requires measuring energies of multi-hadron states.

# Resonance energies and widths

- Above inelastic threshold, even less is known precisely.

- Resonant states will appear as “avoided level crossings” in the spectrum.
- Width can be inferred from the gap at the point where the energy levels get closest.



- Example: two states,  $|\phi\rangle$  and  $|\chi(p)\chi(-p)\rangle$  with  $p = 2\pi/L$ .

# Resonance energies and widths

- Modelling these level crossings can be used to predict the energy and width of the resonance. Extracting these parameters from Monte Carlo data will require a precise scan of the energy of many states (ground-state, first excited, second ...) in a given symmetry channel to be carried out at a number of lattice volumes.

## Requirements for measuring decay widths in QCD

- Light, dynamical quarks (to ensure unitarity)
- Accurate spectroscopy in appropriate channels
- Simulations in multiple box sizes and shapes
- Access to excited states in these channels
- Ability to create multi-hadron states

→ “Next-generation” lattice spectroscopy.

# Conclusions:

- **Good news:**

- 1 The lattice defines field theory without reference to perturbative expansions.
- 2 The lattice regulates quantum fields
- 3 The lattice provides a framework for Monte Carlo simulations on (large) computers.
- 4 Successes include: spectrum, QCD coupling, quark mass determinations, matrix element determinations for CKM physics

- **Bad news:**

- 1 The chiral symmetry of quarks is hard to handle.
  - 2 For most of this to work, need the Wick rotation, so ...
  - 3 Studying scattering is very complicated
  - 4 Studying strong decay is very complicated
  - 5 Still can't run simulations at physical quark masses so need to extrapolate in quark masses.
- There are plenty of other challenges: QCD at finite density, ...